

In-situ co-doping of sputter-deposited TiO₂:WN films for the development of photoanodes intended for visible-light electro-photocatalytic degradation of emerging pollutants

N. Delegan,¹ R. Pandiyan,¹ S. Komtchou², A. Dirany,² P. Drogui,² and M. A. El Khakani^{1,*}

¹*Institut National de la Recherche Scientifique, Centre-Énergie, Matériaux et Télécommunications, 1650 Blvd. Lionel-Boulet, Varennes, QC J3X-1S2, Canada*

²*Institut National de la Recherche Scientifique, Centre-Eau, Terre et Environnement, 490 Rue de la Couronne, QC G1K-9A9, Canada*

*Corresponding author: elkhakani@emt.inrs.ca

Abstract

We report on the magnetron sputtering deposition of *in-situ* codoped TiO₂:WN films intended for electro-photocatalytic (EPC) applications under solar irradiation. By varying the RF-magnetron sputtering deposition parameters, we were able to tune the *in-situ* incorporation of both N and W dopants in the TiO₂ films over a wide concentration range (i.e. 0-9at.% for N and 0-3at.% for W). X-ray photoelectron spectroscopy analysis revealed that both dopants are mostly of a substitutional nature. The analysis of the UV-Vis transmission spectra of the films confirmed that the optical bandgap of both TiO₂:N and TiO₂:WN films can be significantly narrowed (from 3.2eV for undoped-TiO₂ down to ~2.3eV for the doped ones), by tuning their dopant concentrations. We were thus able to pinpoint an optimal window for both dopants (N and W) where the TiO₂:WN films exhibit the narrowest bandgap. Moreover, the optimal codoping conditions greatly reduce the recombination defect state density compared to the monodoped TiO₂:N films. These electronically passivated TiO₂:WN films are shown to be highly effective for the EPC degradation of atrazine (pesticide pollutant) under sunlight irradiation (93% atrazine degraded after only 30 min of EPC treatment). Indeed, the optimally codoped TiO₂:WN photoanodes were found to be more efficient than both the undoped-TiO₂ and equally-photosensitized TiO₂:N photoanodes (by ~70% and ~25%, respectively), under AM1.5 irradiation.

Introduction

2 Since the pioneering work of Fujishima and Honda in 1972 on titanium dioxide (TiO_2) water
3 splitting¹, the research on the photocatalytic properties of TiO_2 and its varied forms has continued
4 to attract great attention. In fact, TiO_2 offers a unique combination of unpaired features among which
5 are its wide bandgap and high chemical stability¹⁻⁵ in addition to its availability and affordability.
6 All these facts have made TiO_2 the candidate of choice for applications in solar energy
7 conversion^{1,2,5} and photo-degradation of contaminants.^{6,7}
8 Fundamentally, TiO_2 is an n-type semiconductor with a wide intrinsic bandgap (E_g) of 3.0 eV and
9 3.2 eV for rutile and anatase polymorphs, respectively. This inherently restricts the photoactivity of
10 titania based devices to the **ultra-violet (UV)** range (which represents only ~4% of the irradiance of
11 the sunlight) and limits to a certain extent its implementation in large scale industrial photo-
12 processes.⁸ It is thus understandable that substantial research effort is aimed at extending the
13 photosensitivity of TiO_2 to the visible light range (which represents ~40% of solar irradiance)
14 through the narrowing of its bandgap via varied approaches including self-doping,⁹ dye-
15 sensitization,^{2,4} cationic doping,^{4,7} and anionic doping.^{3,4,6,7} For the latter approach, substitutional
16 **nitrogen (N) doping** has been shown to be very effective, as it leads to the formation of discrete N_{2p}
17 acceptor states within the E_g above the valence band maximum (VBM) reducing thereby the
18 bandgap energy from 3.2 eV (anatase) to ~2.3 eV (N-doped anatase).^{4,6,7} In this context, we have
19 demonstrated in a previous study⁷ that the *in-situ* N-doping induced E_g narrowing of sputtered TiO_2
20 films was directly correlated with the significant increase of their **electro-photocatalytic (EPC)**
21 capacity to degrade the chlortetracycline (CTC) contaminant under sunlight exposure. In fact, N
22 doping is a more complex process which induces different energy states within the semiconductor's
23 forbidden gap.¹⁰⁻¹⁷ Some of these energy levels act as acceptor states within the E_g and therefore

effectively reduce the energy required for photoexcitation, while others act as charge traps and/or recombination centers hindering thereby the photocharge transport. Notably, the introduction of N into the TiO₂ lattice has been shown to lead to two types of these energy states associated with the dopant being into substitutional or interstitial sites. The substitutional doping is associated with the formation of N_{2p} acceptor levels just above the O_{2p} levels, whereas the interstitial one is related with isolated N-O type deep gap states within the E_g that may act as photocharge recombination centers.¹⁰⁻¹⁵ In addition, the defects induced by the presence of both dopants are usually compensated for by the oxide structure via the formation of oxygen vacancies (V_O), which in turn act as potent recombination centers located beneath the conduction band minimum (CBM).^{14,16} In sum, both experimental and theoretical studies¹⁰⁻¹⁷ indicate that mono-doping (with N for example) while being very effective in narrowing the bandgap, it concomitantly creates recombination centers that tend to reduce the extracted photogenerated current.^{4,7} This limits the full exploitation of the photocatalytic efficiency gain from the absorption edge red-shift.⁵ To counteract this limitation, density of states (DOS) models¹⁸⁻²⁴ have showed that it is possible to minimise the undesired side-effects arising from N-doping by achieving an electronically passivated acceptor-donor co-doping approach. The latter would suppress the charge defects acting as recombination centers, while maintaining an enhanced visible light absorption.^{5,18-21,23,24} By theoretically investigating different dopant couples,^{21,23} tungsten (W) and nitrogen (N) were identified as the most interesting pair that would narrow the E_g while minimizing the recombination of photogenerated excitons through electronic passivation of uncompensated charges.^{18-20,22} Recent chemical synthesis approaches have attempted simultaneous W and N codoping (WN-codoping) of TiO₂ leading to substantial band gap narrowing.²⁵⁻³⁰ However, these studies showed mainly interstitial type of N-doping and the presence of other contaminants such as amino groups and carbonaceous species (inherent to precursors used in chemical methods). So far, the observed increase in EPC activity was mainly associated with the

1 effective bandgap reduction and surface area increase^{25,28,31} with almost no focus and/or
2 optimisation of the electronic passivation induced by the co-doping approach. This calls for a more
3 systematic experimental study of the structural, electronic, and chemical effects resulting from W
4 and N codoping of TiO₂, over the largest dopant range accessible. The objective being to achieve
5 substitutional doping and identify the optimal doping concentrations leading to both narrowest
6 bandgap and a minimum of photocharges trapping and/or recombination (i.e. longer lifetimes of the
7 photogenerated charges). Practically, an effective co-doping of the TiO₂:WN photoanodes will
8 translate into higher EPC photocurrent, which in turn will degrade more efficiently a given
9 contaminant in water.

10 In this paper, we report on the *in-situ* co-doping of TiO₂ films by both W and N atoms by using a
11 reactive magnetron-sputtering process. By concomitantly adjusting the N₂-to-Ar gas flow ratio and
12 the W-to-TiO₂ target power ratio in the RF-sputtering plasma source, we were able to control the W
13 and N incorporation rates into the films, over the (0 - 3) at.% and (0 - 9) at.% concentration ranges,
14 respectively. The effects of both mono-doping (N and W taken separately), and **WN-codoping**
15 approaches on the structural, optical, and chemical properties of the sputtered TiO₂ films have been
16 systematically investigated. By achieving detailed experimental analyses of the energy band
17 structure of the doped TiO₂ films, we have recently demonstrated that optimal W and N doping
18 contents correspond not only to the narrowest E_g value of ~2.2 eV, but also to favorable Fermi level
19 (E_F), CBM, and VBM band positions.³² In the present paper, the codoped TiO₂:WN films are shown
20 to not only possess favorable optoelectronic properties but most interestingly to exhibit also densities
21 of oxygen vacancies and photocharge recombination centers lower than in their mono-doped
22 counterparts, confirming thereby the achievement of an electronically passivated acceptor-donor co-
23 doping approach. To assess the practical advantages of such a co-doping approach, the optimised
24 TiO₂:WN films were integrated into photo-anodes for the EPC degradation of atrazine (a pesticide,

1 and residual contaminant in water) under simulated AM1.5 solar light, and found to deliver superior
2 EPC performance their mono-doped photoanode counterparts.

3

4 **Experimental**

5 All the films were deposited from the sputtering of a TiO₂ target by using an RF (13.56 MHz)
6 magnetron gun operating at a constant power density of 8.8 W/cm². The *in-situ* N-doping of the
7 films was achieved through the introduction of N₂ in conjunction with the Ar sputtering gas, whereas
8 W-doping was obtained through the co-sputtering of both TiO₂ and W targets (3''-diam. and 99.99%
9 purity). For the WN-codoping, both N₂ gas and W target co-sputtering were combined with the main
10 TiO₂ target sputtering. Prior to deposition, the chamber was cryopumped to a base pressure of 2 ×
11 10⁻⁸ Torr. Thereafter, high purity Ar (99.999 %) and N₂ (99.995 %) gases were introduced into the
12 deposition chamber. High purity O₂ (99.995%) was also used for the deposition of the reference
13 TiO₂. The gas flow rates were monitored to maintain a constant pressure of 1.4 mTorr inside the
14 chamber during the sputter-deposition process. In order to control the N incorporation in the
15 TiO₂:WN and TiO₂:N films, the relative nitrogen mass flow rate ratio R_{N2} (i.e. [N₂]/([N₂]+[Ar])) was
16 varied between 0 and 15%. The W incorporation was controlled by varying the W target power
17 density (W_{power}) from 50 mW/cm² to 250 mW/cm². The TiO₂:WN films were simultaneously
18 deposited onto different substrates, including: (i) undoped and double side-polished Si(100), (ii)
19 quartz, and (iii) deployed (10 cm x 11 cm) Ti-grids. The substrates were mounted on a holder located
20 off-axis at a distance of 20 cm from the target, and heated during deposition by a quartz lamp heater
21 to an actual on-substrate temperature of ~470 °C. Prior to film deposition, the target was
22 systematically sputter-cleaned with Ar ions for ~15 min with the shutters closed. The thickness of
23 the TiO₂:WN films was *in-situ* monitored by means of a calibrated quartz-crystal microbalance and
24 *ex-situ* measured through cross-section scanning electron microscopy (SEM) observations by means

1 of a Jeol JSM-6300F SEM system. No post-acceleration voltage was intentionally applied to the
2 substrates during the sputter-deposition process (they were nonetheless subjected to a built-in
3 plasma sheath bias of ~ -14 V during their growth). The atomic composition and chemical bonding
4 states of the films were systematically investigated by means of X-ray photoelectron spectroscopy
5 (XPS). The XPS spectra were collected by using the ESCALAB 220i-XL spectrophotometer,
6 equipped with an $AlK\alpha$ (1486.6 eV) twin-anode source, after a systematic *in-situ* surface cleaning
7 by means of 5 keV Ar^+ ion sputtering gun. The crystalline structure of the films on quartz substrates
8 was characterized by means of a PANalytical X-Pert Pro X-ray diffractometer (XRD) system
9 using $CuK\alpha$ radiation of 1.5418 Å. The **ultra-violet and visible (UV-Vis)** transmittance and
10 reflectance spectra of the $TiO_2:WN$ films were systematically measured by using an Agilent/Varian
11 Cary 5000 system. Thus, their E_g values were derived from UV-Vis Tauc plot extrapolation. The
12 EPC properties of the $TiO_2:WN$ films deposited onto deployed titanium-grid substrates (used as
13 photoanodes) were evaluated in an EPC reactor of which details are described elsewhere³³.

15 **Results and Discussion**

16 The structural and morphological features arising from the concomitant incorporation of N and W
17 atoms into TiO_2 were studied via cross-section SEM observations. We have previously showed that
18 the introduction of N_2 in the deposition chamber leads to an energetic competition between the the
19 dissociation of N_2 molecules and the production ions that contribute to the sputtering process itself.⁷
20 Here, we focus on the concomitant interplay between W and N sources and their effects on the
21 deposited films. Figure 1(a) presents the deposition rate as a function of W_{power} (bottom x-axis; with
22 R_{N_2} kept constant at 6%), and R_{N_2} (top x-axis; with W_{power} kept constant at 132 mW/cm²). We note
23 that N_2 introduction creates a competitive balance between the energy consumption for N_2 molecule
24 dissociation and sputtering ion production with an optimal deposition point at around 7% R_{N_2} . It is

1 expected that this optimal point should shift to higher R_{N_2} as more power is put in the system.
2 Additionally, we note that increasing the power imposed on the W target while maintaining R_{N_2}
3 constant linearly increased the deposition rate. This power dependence of the **deposition rate** is
4 typical for metallic targets in a sputtering process.³⁴ The morphology of the sputter-deposited
5 $TiO_2:WN$ films showed no significant differences as a function of deposition parameters. From that
6 perspective, all the sputter-deposited films were found to be dense, smooth, and exhibit a uniform
7 apparent morphology regardless of W_{power} or R_{N_2} . A typical cross sectional SEM image of a
8 $TiO_2:WN$ film can be seen in Figure 1(b), the deposition parameters for this film were the crossover
9 point from Figure 1(a): $W_{power} = 132 \text{ mW/cm}^2$ and $R_{N_2} = 6\%$. The thicknesses of the films
10 investigated here are in the 200-400 nm range, depending on the deposition conditions. The exact
11 film thickness will be used for each corresponding film for the determination of its optical absorption
12 coefficient.

13 The effectiveness of our sputter process to incorporate N and W dopants in the $TiO_2:WN$ films was
14 investigated by means of XPS analysis of the N $1s$, Ti $2p$, O $1s$, W $4f$, and W $4p$ core level spectra.
15 The variation of deposition parameters W_{power} and R_{N_2} on the incorporation of each dopant is
16 presented in Figure 2. Figure 2(a) shows the incorporation of W (left y-axis) and N (right y-axis)
17 atoms with varying R_{N_2} while W_{power} is maintained constant at 132 mW/cm^2 . From this Figure 2(a),
18 one can see that the increase of R_{N_2} translates into a progressive increase of N incorporation.
19 However, it is also seen that the increase of N incorporation occurs to the detriment of W
20 incorporation. The increase of N content with R_{N_2} is a direct consequence of an increased amount
21 of N atoms in the plasma mix during the deposition process. This results into more N atoms available
22 for incorporation, which enter in a sort of competition with W atoms (at a fixed power on the W
23 target), lowering thereby their incorporation when R_{N_2} is increased. In contrast, when the R_{N_2} is

1 fixed (at 6% here) and W_{power} is increased, the incorporation of both N and W steadily increase with
2 W_{power} , as shown in Figure 2(b). This behavior can be explained by two facts: (i) the incorporation
3 of W into the TiO_2 lattice is expected to increase the solubility limit of the N dopant.^{5,23} The
4 solubility increase is caused mainly by the introduction of W atoms into the TiO_2 lattice which
5 lowers the formation energy of substitutional N and increases the total free valence of the
6 matrix.^{21,23,24} Secondly, (ii) the increase of W_{power} , while keeping a constant R_{N_2} , leads to more
7 sputtered W atoms and concomitantly provides additional energy for the dissociation of N_2
8 molecules into N atoms. Figures 2 (a) and (b) highlight the ability of our process to precisely control
9 the incorporation levels of each dopant in the sputtered TiO_2 films through the appropriate tuning of
10 both W_{power} and R_{N_2} . Finally, the fundamental characteristics of the XPS binding energy (BE) spectra
11 were found to be very similar within the investigated concentration ranges (i.e. 0 to ~3 at.% and 0
12 to ~9 at.% for W and N contents, respectively). Indicating that for these doping levels, the chemical
13 bonding of these films may vary in content (depending on the doping levels), but their nature remains
14 the same.

15 Figure 3 shows the high-resolution BE spectra of the O $1s$, Ti $2p$, N $1s$, and W $4f$ (from left to right
16 respectively) core levels for different films, namely TiO_2 , $\text{TiO}_2\text{:N}$, $\text{TiO}_2\text{:W}$, and $\text{TiO}_2\text{:WN}$ films
17 (from bottom to top, respectively). These core level spectra were deconvoluted into their different
18 components corresponding to various local electronic environments (oxidation/bonding states). For
19 the TiO_2 films, the Ti $2p_{3/2}$ peak can be decomposed into two main components corresponding to
20 two oxidation states of Ti, namely Ti^{IV} and Ti^{III} at 458.6 eV and 456.9 eV, respectively, with their
21 corresponding Ti $2p_{1/2}$ energy levels found at 5.7 eV and 5.5 eV higher BE, having 1:2 branching
22 ratios.^{7,31,35-37} For subsequent W $4f$ fitting purposes (as the Ti $3p$ energy window overlaps the W $4f$
23 peak energies), Ti $3p$ peaks were deconvoluted to reveal Ti^{IV} and Ti^{III} oxidation state peaks at 37.1

1 eV and 35.7 eV respectively,^{38,39} in full agreement with the Ti $2p$ derived quantifications. The
2 presence of Ti^{III} can be attributed to sub-stoichiometric species formed mainly as a result of
3 preferential sputtering during thin film deposition.^{7,40} Given that Ar^+ ion etch-cleaning is used to
4 eliminate surface carbonaceous contaminants prior to acquisition, the Ti^{III} quantification is assumed
5 to be slightly overestimated. However, this contribution is expected to remain unchanged for all
6 surveyed samples as surface etch-cleaning parameters were kept constant.^{7,36,40} This feature is
7 important as its presence is indicative of V_O in the lattice, a defect associated with high levels of
8 photocharge recombination.^{4,8,41} The Ti $2p$ states can be inferred in the corresponding O $1s$ core
9 levels. The main $O^{(TiIV)}$ peak at 529.9 eV is associated with lattice oxygen in the TiO_2 matrix; we
10 note that sub-stoichiometric lattice oxygen can also be distinguished and identified as $O^{(TiIII)}$ at
11 ~ 0.5 eV higher binding energy.^{35,36} Additionally, we note the presence of a third peak, denoted as
12 $O^{(H)}$ at 530.7 eV and about double the FWHM of the other O $1s$ peaks. Commonly, this peak is
13 attributed to O $1s$ defective oxygen and surface hydroxide species.^{35,36,42} This contribution is
14 expected to remain of similar relative intensity in relation to the main O $1s$ peaks for all films.^{35,36}
15 The XPS derived composition of undoped films, calculated from relative peak intensities, is found
16 to be $\sim TiO_{1.87}$.

17 As nitrogen gas is introduced into the deposition chamber ($TiO_2:N$ films), the appearance of the N
18 $1s$ peak in the film's XPS spectra can be observed with two characteristic components: N^{sub} and
19 N^{inter} at 396.2 eV and 397.2 eV binding energies, respectively. The N^{sub} peak is known to arise from
20 substitutional -Ti-N- type bonds and the N^{inter} peak is associated with interstitial -O-N type bonds
21 within the titania crystal structure.^{4,6,7,11-13} The relative intensity of the peaks confirms that most of
22 the doping is of the substitutional type, associated with the visible light photosensitisation of $TiO_2:N$
23 via the introduction of N $2p$ levels just above the O $2p$ composed VBM.^{4,6,7,11-13} However, the

1 introduction of N into the thin films is accompanied with the appearance of Ti^{II} oxidation state
2 species in the Ti 2*p* BE spectra (Ti^{II} 2*p*_{3/2} at 455.3 eV and the 2*p*_{1/2} component at 5.6 eV higher BE)
3 and Ti 3*p* (Ti^{II} 3*p* at 34.8 eV).^{4,43} This reduced state can also be inferred on the O 1*s* spectra with
4 the appearance of O^(TiII) (~1eV higher BE than the O^(TiIV) peak).^{35,38} These reduced valence states
5 are known to arise from the electronic distortions induced by nitrogen doping, reducing Ti^{IV} species
6 into Ti^{III}, and on a much lesser extent into Ti^{II}. These bonding states are indicative of the lattice
7 compensating for localised deep charge defects (known to arise from interstitial type N doping),
8 mainly via the formation of V_O.^{4,7,8,14,41} If substitutional type of doping is considered, the obtained
9 stoichiometry of the N-doped films is ~TiO_{1.69}:N_{0.12} (total N ~6.3 at.%, with N_{sub} ~4.2 at.%); this
10 relatively low O content indicates a high density of V_O present within the thin films, justifying the
11 significant increase of highly reduced Ti species.

12 The addition of tungsten (W) in the co-sputtering process, via sputtering of a metallic W target
13 concomitantly with the TiO₂ target sputtering, leads to the appearance of two main features in the
14 W 4*f* BE window of the TiO₂:W thin-films (Figure 3). These W 4*f* peaks are due to the W^{VI} and W^{IV}
15 states in the oxide form with the main 4*f*_{7/2} components appearing at 34.5 eV and 32.3 eV,
16 respectively, with the corresponding 4*f*_{5/2} doublets found at 2.2eV higher BE with a 3:4 peak
17 intensity ratio.^{37,44-47} The appearance of the O^(W) peak at 529.2 eV in the O 1*s* spectra can be inferred,
18 as some of oxygen species are associated with W.⁴⁶ The XPS derived composition of these films is
19 ~TiO_{1.84}:W_{0.06} (total W ~2.3 at.%). From these analyses, we note that W doped films present lower
20 density of V_O than N doped films. This aligns with theoretical models predicting that W-doping
21 should reduce the formation of V_O (increase in formation energy) as compared to N-doping.²⁰

22 Due to size constraints, W and N are mostly expected to substitute for Ti and O atoms, respectively
23 (nonetheless, interstitial N is also possible as above discussed), and -W-N- type bonds are

energetically superimposed to -Ti-N- type bonds.^{48,49} Therefore, little effect is expected in the N *1s* and corresponding cation signatures. Thus, it can be confirmed that both the W and N species are incorporated into mostly substitutional locations within the TiO₂ matrix as they present identical chemical energy states as the TiO₂:N^{4,6,7} and TiO₂:W.^{37,44–46} This being said, the XPS derived composition of this codoped film is ~TiO_{1.83}:W_{0.04}N_{0.09} (N ~4.9 at.% and W ~1.2 at.%). This film composition indicates that there is a lack of O in the films, at a similar atomic proportions as in the undoped samples (which were interestingly deposited with additional oxygen during sputtering). This indicates that there are less oxygen vacancies in the codoped films as compared to the N-monodoped ones. This is in accordance with theoretical models that predict that when both dopants (W and N) are introduced simultaneously, there is a significant reduction of electronic defect formation (such as V_O).^{5,18–21,23,24} Table 1 summarises the XPS derived dopant levels of the TiO₂:WN films along with their corresponding thickness, as determined from SEM cross-section observations.

N (at.%)	W (at.%)	Thickness (nm)
0.0	0.0	200
2.4	0.0	260
5.7	0.0	250
0.0	0.6	400
0.0	2.3	380
0.0	3.1	430
2.5	2.4	300
5.9	0.7	310
6.4	1.1	300
7.2	1.6	300
8.9	2.4	410
4.9	1.2	370
8.5	0.9	270
7.7	0.2	270
7.9	0.3	280
8.7	3.6	340
6.4	3.6	270

Table 1: Film thickness (as derived from SEM observations) and total dopant contents (from XPS) for the samples used in the UV-Vis. characterizations.

Our previous work has shown that the N substitutional doping significantly affects the optical properties of the N-doped TiO₂ films, namely narrowing their bandgap.⁷ Hence, the effect of W and N dual doping on the optical properties of the TiO₂:WN films was systematically investigated as a function of the doping scheme (i.e., N, W, and WN). Figure 4(a) shows UV-Vis transmittance spectra of the TiO_{1.87}, TiO_{1.84}:W_{0.06}, TiO_{1.69}:N_{0.12}, and TiO_{1.83}:W_{0.04}N_{0.09} films deposited onto quartz substrates. We note a significant red shift of the absorption band edge of the spectra from about 320nm for undoped TiO₂ films to ~420nm for the TiO_{1.69}:N_{0.12} films, and to ~440 nm for TiO_{1.83}:W_{0.04}N_{0.09} films. In agreement with our previous work, our results showed that any doping content beyond this doping point did not necessarily translate into further red shifting of the absorption edge.⁷ It is to be noted that the W-doped films (TiO_{1.84}:W_{0.06}) showed the same absorption edge as undoped TiO₂ films (no red-shift). This confirms that it is indeed nitrogen that produces energy states responsible for the observed photosensitization of the TiO₂:WN films in the visible. In order to derive a quantitative value of the E_g for these thin films, the corresponding UV-Vis spectra (transmittance and reflectance used to extract the absorption coefficient) were analyzed according to the Tauc formalism,⁴⁵ where the E_g value can be deduced from the x-intercept of the tangential extrapolation of the Tauc plots, as shown in Figure 4(b). We note that indeed, both the TiO_{1.69}:N_{0.12}, and TiO_{1.83}:W_{0.04}N_{0.09} samples show significant optical E_g reduction to around ~2.3 eV, while the TiO_{1.87} and TiO_{1.84}:W_{0.06} films exhibit an E_g of ~3.2 eV, consistent with indirect gap values quoted for titania in the literature.^{4,8} These results confirm that W doping has no significant impact on the absorption edge of the films, while the N doping induced visible photosensitisation is still effective in the co-doping scheme. In order to get a broader picture, the E_g values were derived for numerous TiO₂:WN codoped films with a broader range of doping levels and reported in a sort of a 2D map

of which X and Y axes are N and W doping level, as shown in Figure 5. We note that in the Y-axis (W concentration) that there is a very little to no effect on the E_g variation (varying W by about 3 at.% induces only a small ~ 0.2 eV change in E_g value), mostly through the creation of shallow states below the CB.^{19,50} In contrast, the main determining factor for bandgap narrowing of the TiO₂:WN films (and thereby their photosensitisation to visible light) is their N doping level. Additionally, we noticed that beyond a N-doping level ~ 6 at.% of N, no additional gain in visible light absorption can be obtained. This points up the existence of an optimal nitrogen doping concentration range (5-6 at.%) where the E_g stays at its lowest value of ~ 2.3 eV (a bandgap value that ensures effective visible light photosensitization of TiO₂ films). This observation, along with the theoretical prediction²⁰ that states that the acceptor-donor passivation is most efficient with the N:W ratio close to 2:1 is what guided the selection of our representative TiO₂:WN samples to be used for further analysis. Thus, the following film compositions were investigated: TiO_{1.87}, TiO_{1.84}:W_{0.06}, TiO_{1.69}:N_{0.12}, and TiO_{1.83}:W_{0.04}N_{0.09}. In order to better understand the relationship between EPC performance and dopant induced band-gap narrowing, one has to have access to other physical parameters, such as the Fermi energy, work function, and valence band maximum (VBM)/conduction band minimum (CBM) energy levels (a study that is beyond the scope of the present paper). In fact, the measurement and quantification of the above-mentioned optoelectronic characteristics of the doped TiO₂ films in order to reconstruct the band energy diagrams is a thorough study that has been reported elsewhere.³² We briefly highlight here some of the relevant insights gained from this study³²: first, (i) W-doping keeps most of the E_g features unchanged outside of the introduction of new states below the CBM, (ii) N-doping shrinks the E_g mostly via the introduction of new mid-gap N $2p$ states above the mostly O $2p$ populated VBM while reducing the work function, and (iii) WN-codoping shrinks the E_g mostly via hybridizing N $2p$ states with VBM O $2p$ states while leaving most other features unchanged.³² In summary, while both N-doped and WN-codoped samples show a significant

1 reduction in their E_g , the acceptor-donor codoped samples are expected to outperform monodoped
2 ones due to less midgap states, defect states, and better placements of the VBM energy levels.³²

3 It is worth reiterating that the crystallinity of the films plays a key role in their EPC properties as it
4 determines features such as the bandgap, average crystallite size (mean free path of charges), etc^{4,8}.
5 Therefore, we have used XRD to investigate the crystalline structure of our films as a function of
6 their doping scheme, and the obtained spectra are displayed in Figure 6. The $\text{TiO}_{1.87}$ films are seen
7 to crystallize in the anatase (A) phase with the characteristic (101), (103), (004), (112), (200), (105),
8 and (211) peaks (JCPDS no.: 84-1286); with the main A(101) peak at 25.4° with 0.4° of full width
9 at half maximum (FWHM). The Williamson-Hall size-strain analysis⁵¹ revealed lattice parameters
10 typical of anatase under low tensile strain ($\sim 0.66\%$), as shown in Table 2. This is most likely due to
11 the oxygen vacancies present in the structure, causing some of the Ti atoms to reduce from Ti^{IV} to
12 Ti^{III} (as above-shown by XPS), increasing their atomic radii. When nitrogen is incorporated into the
13 structure, the crystalline quality of the $\text{TiO}_2\text{:N}$ films is significantly reduced as seen in Figure 6
14 (signal to noise ratio and FWHM $\sim 0.7^\circ$ for the A(101)). The tensile strain of the $\text{TiO}_2\text{:N}$ films
15 increased to $\sim 2.75\%$, explaining the observed relative deterioration of their crystalline quality. This
16 strain exists due to a couple of reasons: (i) N-doping reduces the formation energy of V_O facilitating
17 sub-stoichiometric oxide formation,^{4,23,24,52,53} and (ii) interstitial N-doping will impose a certain
18 stretching of the lattice.⁵⁴ In the case of W doping, the crystallinity of the $\text{TiO}_2\text{:W}$ films is also seen
19 to be greatly affected. Given that W^{VI} and Ti^{IV} are of similar ionic radii (around 0.600 \AA and 0.605
20 \AA), substitutional W^{VI} doping should not affect the TiO_2 lattice structure significantly. However, as
21 shown by XPS analysis in Figure 3 we note the presence of W^{IV} species in the TiO_2 structure, this
22 reduced state is of higher atomic radii (0.660 \AA) and is associated with the distortion of the crystal
23 lattice to higher 2θ values.³⁷ Additionally, W^{IV} species (WO_2 type crystal) naturally crystallize in a

1 rutile structure.⁵⁵ In fact, this and the W doping generated strain (as high as ~4.06%) in the lattice
 2 induce the TiO₂:W films to transition into the rutile polymorph.^{7,37} In Figure 6 the rutile phase peaks
 3 are seen to be predominant and very broad (with a FWHM as large as ~2.1° for the R(110) peak).
 4 The less intense anatase peak (A(101)) is also found to be very broad with a FWHM of ~1.9°).
 5 Interestingly, in the case of **WN-codoping**, the anatase phase is recovered with a high crystalline
 6 quality (the FWHM of the anatase (101) peak FWHM is of ~0.5°) and a moderate tensile strain
 7 (estimated to ~1.9%). The crystallite size (of ~30 nm) of these codoped films was found to be very
 8 comparable to that of undoped TiO_{1.87} films. Moreover, one can note that the strain in the **WN-**
 9 **codoped** films is lower as compared to their monodoped counterparts (i.e. W-doped and N-doped;
 10 see Table 2). **This recovery of the anatase phase seems to be a direct result of local dopant charge**
 11 **passivation, leading to a reduction of structural defects that arise from uncompensated lattice**
 12 **charges. Indeed, substitutional W and N dopants have been reported to result in a charge-passivated**
 13 **lattice with lower V_O formation.^{21–24}**

Sample	TiO ₂	TiO _{1.89}	TiO _{1.84} :W _{0.06}	TiO _{1.69} :N _{0.12}	TiO _{1.83} :W _{0.04} N _{0.09}
a (nm)	0.378	0.378	N/A	0.381	0.384
b (nm)	0.378	0.378	N/A	0.381	0.384
c (nm)	0.950	0.944	N/A	0.966	0.934
c/a	0.251	0.250	N/A	0.254	0.243
Vol. (nm ³)	0.136	0.135	N/A	0.140	0.138
Crys. Size (nm)	N/A	27.68	6.40	7.00	32.93
Internal strain (η)	0.00%	0.66%	4.10%	2.75%	1.90%

14 Table 2: Williamson-Hall calculated lattice parameters based on the XRD spectra of Figure 4 for
 15 the various doping schemes alongside theoretical TiO₂ lattice values.

16 Finally, in order to assess the electro-photocatalytic properties of our doped TiO₂ films, while taking
 17 advantage of their visible light photosensitisation and the acceptor-donor passivated codopants, the
 18 codoped TiO_{1.83}:W_{0.04}N_{0.09} films along with TiO_{1.87}, TiO_{1.69}:N_{0.12}, and TiO_{1.84}:W_{0.06} films were

1 deposited onto both sides of deployed Ti-grid substrates (10 cm x 11 cm) and used as photoanodes.
2 Those photoanodes were integrated into a home built electro-photocatalytic (EPC) reactor (of which
3 details can be found elsewhere³³). The EPC performance of our photoanodes was assessed towards
4 the degradation of a real emerging contaminant that is atrazine. Atrazine is a pollutant that is widely
5 used as a pesticide in agriculture. For the purpose of the present study, synthetic solutions containing
6 atrazine concentrations of 60 $\mu\text{g/L}$ (60 ppb, as these are the typical concentrations of heavily polluted
7 areas) were prepared and treated in our EPC reactor under AM 1.5G solar simulator while
8 continuously circulating the solution during the treatment time which was of up to 2 hours. The
9 reaction kinetics for the four different types of photoanodes are shown in Figure 7(a). We note that
10 while $\text{TiO}_{1.87}$ is found to degrade atrazine, the degradation reaction is relatively slow (after 1 hour
11 of treatment, ~ 20 ppb of atrazine I still present in the solution). In the case of $\text{TiO}_{1.69}\text{:N}_{0.12}$
12 photoanodes, the degradation efficiency is seen to improve due to the photosensitisation effect of
13 nitrogen doping; allowing the use of a higher portion of visible light photons for the generation of
14 excitons to be used in the photocatalytic process.^{7,33} Surprisingly, $\text{TiO}_{1.84}\text{:W}_{0.06}$ photoanodes were
15 found to perform similarly (within the statistical error margin) to the photosensitized $\text{TiO}_{1.69}\text{:N}_{0.12}$,
16 even if their bandgap is much larger than that of N-doped TiO_2 . This effect is mainly attributed to
17 the W induced increase of the photocharge lifetimes of the $\text{TiO}_2\text{:W}$ films, as compared to the N-
18 doped photoanodes.^{19,50,56} In fact, a more systematic study of the effect of the doping scheme on the
19 photocharge lifetimes has been carried out and reported elsewhere.⁵⁶ Important to underline here is
20 that the W dopant was found to increase the formation energy of V_O defects, and to locally passivate
21 charge disparities when co-introduced with N.⁵⁶ These effects have led to WN-codoped samples
22 exhibiting a very good visible-light sensitivity while having photocharge lifetimes up to three times
23 longer than those measured in their N-doped TiO_2 counterparts (photocharge trapping decay
24 constant of 90 ± 7 ns for $\text{TiO}_2\text{:WN}$ versus 30 ± 10 μs for $\text{TiO}_2\text{:N}$).⁵⁶ This corroborates well with the

1 fact that the codoped TiO₂:WN photoanodes are the most effective ones for atrazine degradation, as
2 can be clearly seen in Figure 7(b). Indeed, after just 30 min of treatment, the TiO₂:WN photoanodes
3 degraded 94% of the atrazine present in the solution (this represents 24% and 71% more than the N-
4 doped and the undoped TiO₂ photoanodes, respectively). Those results clearly confirm the benefit
5 of the TiO₂ photoanode WN-codoping for the EPC degradation of pollutants. In fact, nitrogen
6 doping allows for photosensitivity of the photoanodes in the visible domain, permitting thereby the
7 additional use of visible light photons for photocatalysis. However, part of this advantage is
8 counteracted by N doping induced deep states in the E_g that act as recombination centers trapping
9 the photogenerated charges. This drawback, in turn, can be minimised via the electronically
10 passivated WN-codoping approach, allowing for the recovery of a significant increase in EPC
11 performance,^{21,23,24} as exhibited here by our TiO₂:WN photoanodes.

12 Conclusion

13 In summary, we were able to develop and apply a reliable RF magnetron sputtering process for the
14 fabrication on in-situ doped and codoped TiO₂ films with different W and N doping concentrations,
15 in the ranging 0-3 at.% and 0-9 at.% ranges, respectively. All the deposited films have shown a
16 dense, smooth, and uniform apparent morphology, regardless of their doping conditions. By
17 adjusting the relevant operating parameters, namely W_{power} and R_{N2}, we were able to fine-tune the
18 incorporation of W and N dopants into the TiO₂ films. The XPS analyses have revealed that both
19 the W and N species are mostly incorporated into substitutional doping locations in the TiO₂ matrix.
20 The N-doping was shown to increase significantly the density of V_O defects (TiO_{1.69}:N_{0.12}
21 composition as compared to the undoped TiO_{1.87}). In contrast, the (WN) codoping was found to
22 minimise the formation of V_O (i.e.; TiO_{1.83}:W_{0.04}N_{0.09}). From optoelectronic viewpoint, both N-
23 doping and WN-codoping were found to induce a significant red shift of the absorption edge, leading

1 to a bandgap narrowing by ~ 0.9 eV. Thus, the bandgap of both $\text{TiO}_2:\text{N}$ and $\text{TiO}_2:\text{WN}$ films was
2 found to be in the visible range with a value as low as 2.3 eV. N-doping was pinpointed as the most
3 determining factor for bandgap narrowing. An N-doping concentration of ~ 5.5 at.% was identified
4 as the optimal content for the narrowest bandgap. Higher N-doping contents do not translate into
5 additional narrowing of the bandgap. On the other hand, XRD results have shown a significant
6 decrease in crystalline quality of the films and an increase in their lattice strain in the case of
7 monodoping. This was particularly observed in the case of W-doped films, where the relatively
8 strong tensile strain has led to a transition from anatase to rutile phase. This drawback of
9 monodoping was significantly reduced when the codoping scheme was adopted, as the lattice strain
10 was found to diminish significantly along with the recovery of high crystallinity (of the anatase
11 phase) of the codoped films. Finally, our EPC experiments have shown that $\text{TiO}_2:\text{WN}$ codoped
12 photoanodes are the most effective for atrazine degradation (93% of atrazine was degraded within
13 the first 30 min of treatment time under AM1.5G solar radiation). Indeed, the codoped $\text{TiO}_2:\text{WN}$
14 photoanodes have performed 71% better than the undoped- TiO_2 and 24% better than the N-doped
15 $\text{TiO}_2:\text{N}$ ones, during the first 30 min of treatment. This emphasizes the benefit of using WN-codoped
16 TiO_2 for EPC applications through the enhanced harvesting of visible light (as evidenced by
17 improved EPC performance) in comparison to N-doped photoanodes. The developed $\text{TiO}_2:\text{WN}$
18 photoanodes definitely open up new prospects for the degradation of emerging contaminants in
19 water while using the direct sunlight radiation.

20 Acknowledgements

21 The authors would like to acknowledge the financial support from NSERC (the Natural Science and
22 Engineering Research Council of Canada), and the FRQNT (Le Fonds de Recherche du Québec-
23 Nature et Technologies) through its strategic Network "Plasma-Québec".

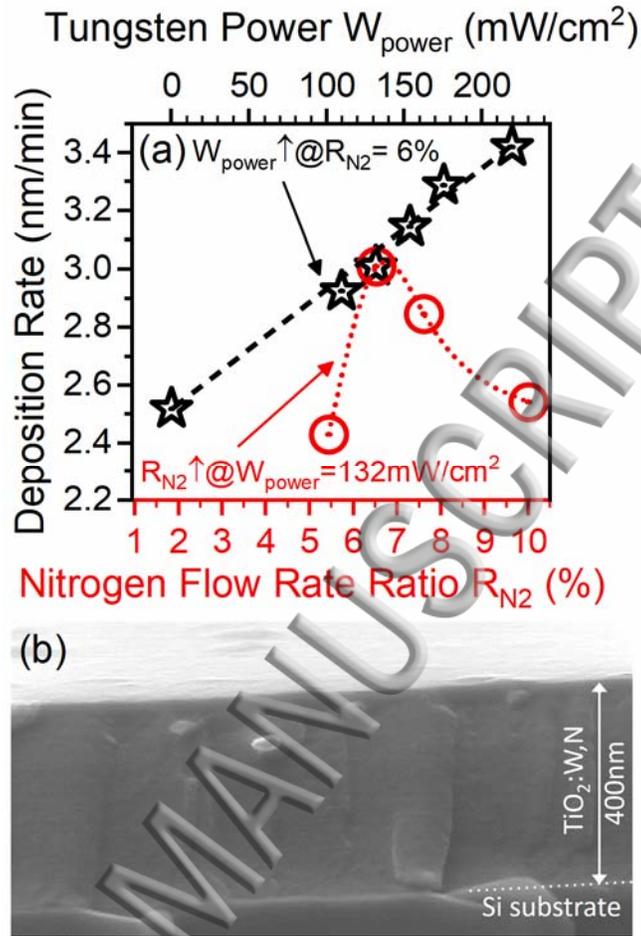
References

- 1 A. FUJISHIMA and K. HONDA, Nature **238**, 37 (1972).
- 2 ¹
- 3 ² B. O'Regan and M. Grätzel, Nature **353**, 737 (1991).
- 4 ³ S.U.M. Khan, M. Al-Shahry, and W.B. Ingler, Science **297**, 2243 (2002).
- 5 ⁴ A. V. Emeline, V.N. Kuznetsov, V.K. Rybchuk, and N. Serpone, Int. J. Photoenergy **2008**, 1
- 6 (2008).
- 7 ⁵ I.S. Cho, C.H. Lee, Y. Feng, M. Logar, P.M. Rao, L. Cai, D.R. Kim, R. Sinclair, and X. Zheng,
- 8 Nat. Commun. **4**, 1723 (2013).
- 9 ⁶ R. Asahi, Science (80-.). **293**, 269 (2001).
- 10 ⁷ N. Deegan, R. Dagherir, P. Drogui, and M.A. El Khakani, J. Appl. Phys. **116**, 153510 (2014).
- 11 ⁸ A. Fujishima, X. Zhang, and D. Tryk, Surf. Sci. Rep. **63**, 515 (2008).
- 12 ⁹ D.O. Scanlon, C.W. Dunnill, J. Buckeridge, S. a Shevlin, A.J. Logsdail, S.M. Woodley, C.R. a
- 13 Catlow, M.J. Powell, R.G. Palgrave, I.P. Parkin, G.W. Watson, T.W. Keal, P. Sherwood, A.
- 14 Walsh, and A. a Sokol, Nat. Mater. **12**, 798 (2013).
- 15 ¹⁰ C. Di Valentin, G. Pacchioni, A. Selloni, S. Livraghi, and E. Giamello, J. Phys. Chem. B **109**,
- 16 11414 (2005).
- 17 ¹¹ C. Di Valentin, G. Pacchioni, and A. Selloni, Phys. Rev. B **70**, 85116 (2004).
- 18 ¹² J. Lynch, C. Giannini, J.K. Cooper, A. Loiudice, I.D. Sharp, and R. Buonsanti, J. Phys. Chem. C
- 19 **119**, 7443 (2015).
- 20 ¹³ F. Peng, L. Cai, H. Yu, H. Wang, and J. Yang, J. Solid State Chem. **181**, 130 (2008).
- 21 ¹⁴ G.R. Torres, T. Lindgren, J. Lu, C.-G. Granqvist, and S.-E. Lindquist, J. Phys. Chem. B **108**,

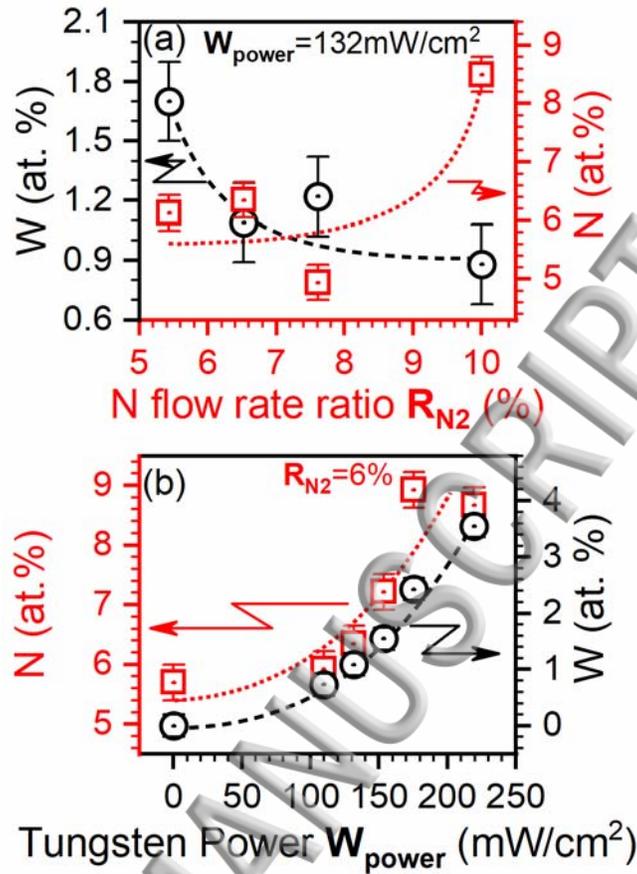
- 1 5995 (2004).
- 2 ¹⁵ M. D'Arienzo, N. Siedl, A. Sternig, R. Scotti, F. Morazzoni, J. Bernardi, and O. Diwald, J.
3 Phys. Chem. C **114**, 18067 (2010).
- 4 ¹⁶ B. Liu, L. Wen, and X. Zhao, Mater. Chem. Phys. **106**, 350 (2007).
- 5 ¹⁷ R. Katoh, A. Furube, K. Yamanaka, and T. Morikawa, J. Phys. Chem. Lett. **1**, 3261 (2010).
- 6 ¹⁸ W.-J. Yin, H. Tang, S.-H. Wei, M.M. Al-Jassim, J. Turner, and Y. Yan, Phys. Rev. B **82**, 45106
7 (2010).
- 8 ¹⁹ V. Çelik and E. Mete, Phys. Rev. B **86**, 205112 (2012).
- 9 ²⁰ M. Li, J. Zhang, and Y. Zhang, Chem. Phys. Lett. **527**, 63 (2012).
- 10 ²¹ R. Long and N.J. English, Chem. Mater. **22**, 1616 (2010).
- 11 ²² R. Long and N.J. English, Appl. Phys. Lett. **94**, 132102 (2009).
- 12 ²³ D. Wang, Y. Zou, S. Wen, and D. Fan, Appl. Phys. Lett. **95**, 12106 (2009).
- 13 ²⁴ Y. Gai, J. Li, S.-S. Li, J.-B. Xia, and S.-H. Wei, Phys. Rev. Lett. **102**, 36402 (2009).
- 14 ²⁵ X. Cui, S. Rong, Y. Cao, Y. Yin, S. Li, and M. Li, Appl. Phys. A **113**, 47 (2013).
- 15 ²⁶ S.S. Thind, G. Wu, M. Tian, and A. Chen, Nanotechnology **23**, 475706 (2012).
- 16 ²⁷ S.S. Thind, G. Wu, and A. Chen, Appl. Catal. B Environ. **111–112**, 38 (2012).
- 17 ²⁸ A. Kubacka, B. Bachiller-Baeza, G. Colón, and M. Fernández-García, Appl. Catal. B Environ.
18 **93**, 274 (2010).
- 19 ²⁹ A. Folli, J.Z. Bloh, E.-P. Beukes, R.F. Howe, and D.E. Macphee, J. Phys. Chem. C **117**, 22149
20 (2013).
- 21 ³⁰ J.Z. Bloh, A. Folli, and D.E. Macphee, J. Phys. Chem. C **118**, 21281 (2014).

- 1 ³¹ S. S. Thind, G. Wu, M. Tian, and A. Chen, *Nanotechnology* **23**, 475706 (2012).
- 2 ³² R. Pandiyan, N. Delegan, A. Dirany, P. Drogui, and M.A. El Khakani, *J. Phys. Chem. C* **120**,
- 3 631 (2016).
- 4 ³³ R. Daghrir, P. Drogui, I. Ka, and M.A. El Khakani, *J. Hazard. Mater.* **199–200**, 15 (2012).
- 5 ³⁴ L.I. Maissel, R. Glang, and P.P. Budenstein, *J. Electrochem. Soc.* **118**, 114C (1971).
- 6 ³⁵ M.C. Biesinger, B.P. Payne, A.P. Grosvenor, L.W.M. Lau, A.R. Gerson, and R.S.C. Smart,
- 7 *Appl. Surf. Sci.* **257**, 887 (2010).
- 8 ³⁶ M.J. Jackman, A.G. Thomas, and C. Muryn, *J. Phys. Chem. C* **119**, 13682 (2015).
- 9 ³⁷ S. Sathasivam, D.S. Bhachu, Y. Lu, N. Chadwick, S. a. Althabaiti, A.O. Alyoubi, S.N. Basahel,
- 10 C.J. Carmalt, and I.P. Parkin, *Sci. Rep.* **5**, 10952 (2015).
- 11 ³⁸ F. Werfel and O. Brümmer, *Phys. Scr.* **28**, 92 (1983).
- 12 ³⁹ J. Riga, C. Tenret-Noël, J.J. Pireaux, R. Caudano, J.J. Verbist, and Y. Gobillon, *Phys. Scr.* **16**,
- 13 351 (1977).
- 14 ⁴⁰ D. Brassard, M. a. El Khakani, and L. Ouellet, *J. Appl. Phys.* **102**, 34106 (2007).
- 15 ⁴¹ A.K. Rumaiz, J.C. Woicik, E. Cockayne, H.Y. Lin, G.H. Jaffari, and S.I. Shah, *Appl. Phys. Lett.*
- 16 **95**, 262111 (2009).
- 17 ⁴² M. Kitano, K. Funatsu, M. Matsuoka, M. Ueshima, and M. Anpo, *J. Phys. Chem. B* **110**, 25266
- 18 (2006).
- 19 ⁴³ M.-S. Wong, H. Pang Chou, and T.-S. Yang, *Thin Solid Films* **494**, 244 (2006).
- 20 ⁴⁴ T. Mishra, M. Mahato, N. Aman, J.N. Patel, and R.K. Sahu, *Catal. Sci. Technol.* **1**, 609 (2011).
- 21 ⁴⁵ J. Li, J. Xu, W.L. Dai, H. Li, and K. Fan, *Appl. Catal. B Environ.* **82**, 233 (2008).

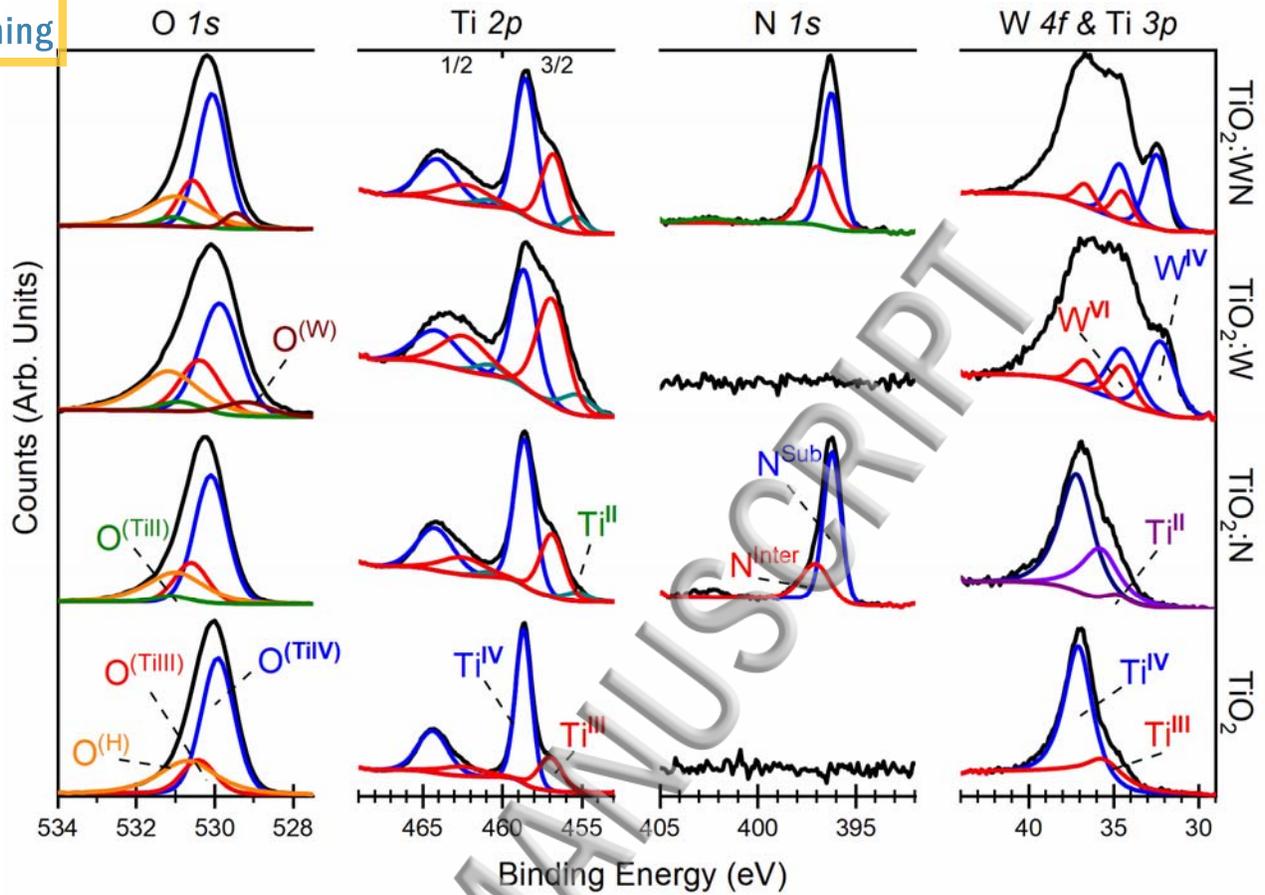
- 1 ⁴⁶ K. T. Ng and D.M. Hercules, J. Phys. Chem. **80**, 2094 (1976).
- 2 ⁴⁷ M. Sun, N. Xu, Y.W. Cao, J.N. Yao, and E.G. Wang, J. Mater. Res. **15**, 927 (2000).
- 3 ⁴⁸ Y.-C. Nah, I. Paramasivam, R. Hahn, N.K. Shrestha, and P. Schmuki, Nanotechnology **21**,
- 4 105704 (2010).
- 5 ⁴⁹ Y. Zhao, W. Hu, Y. Xia, E. Smith, Y. Zhu, C. Dunnill, and D. Gregory, J. Mater. Chem. (2007).
- 6 ⁵⁰ X.. Li, F.. Li, C.. Yang, and W.. Ge, J. Photochem. Photobiol. A Chem. **141**, 209 (2001).
- 7 ⁵¹ G.. Williamson and W.. Hall, Acta Metall. **1**, 22 (1953).
- 8 ⁵² A. Kumar and T. Mohanty, J. Phys. Chem. C **118**, 7130 (2014).
- 9 ⁵³ L. Wan, J.F. Li, J.Y. Feng, W. Sun, and Z.Q. Mao, Mater. Sci. Eng. B **139**, 216 (2007).
- 10 ⁵⁴ Y. Shen, T. Xiong, T. Li, and K. Yang, Appl. Catal. B Environ. **83**, 177 (2008).
- 11 ⁵⁵ F.J. Wong and S. Ramanathan, J. Mater. Res. **28**, 2555 (2013).
- 12 ⁵⁶ N. Deegan, R. Pandiyan, S. Johnston, A. Dirany, S. Komtchou, P. Drogui, and M.A. El
- 13 Khakani, J. Phys. Chem. C **7b11266** (2018).
- 14
- 15



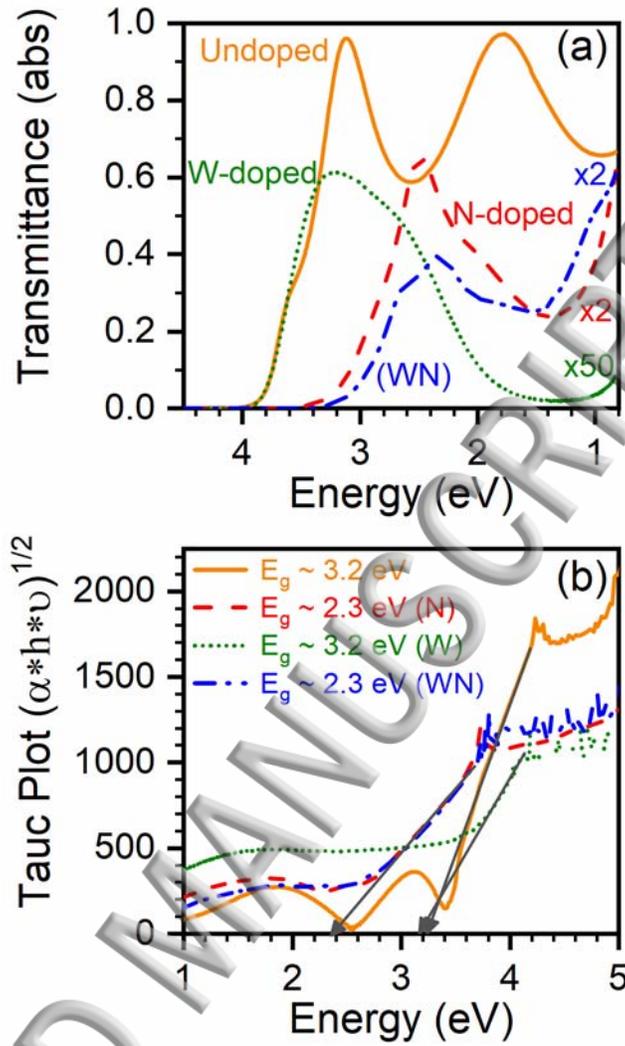
2
3 Figure 1: (a) variation of deposition rate at constant R_{N_2} (black dash-dot line) with varying W_{power}
4 and the variation with constant W_{power} with varying R_{N_2} (red dash line); (b) SEM micrograph of a
5 $W_{power} = 132 \text{ mW/cm}^2$ and $R_{N_2} = 6\%$ deposited $\text{TiO}_2:\text{WN}$ film.



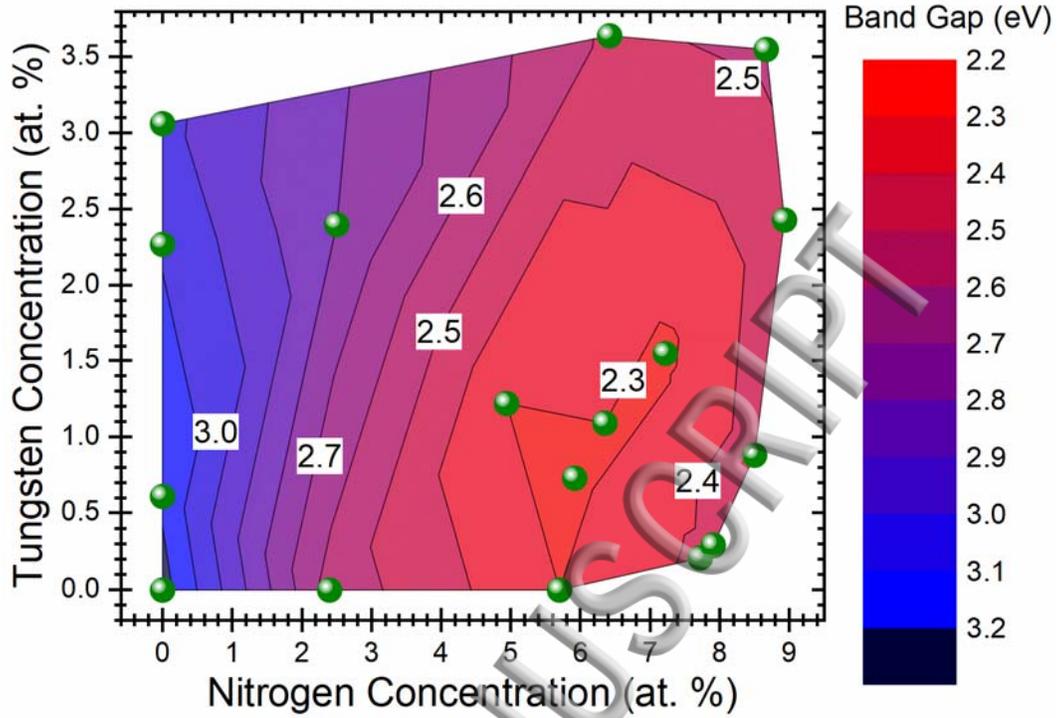
1
2 Figure 2: (a) W and N incorporation rate (left and right hand y-axis respectively) at constant
3 W_{power} with varying R_{N_2} , (b) N and W incorporation (left and right hand y-axis respectively) at
4 constant R_{N_2} with varying W_{power} .



1
2 Figure 3: High-resolution XPS spectra of the O 1s (a), Ti 2p (b), N 1s (c), and W 4f core levels
3 (from left-to-right) of the TiO₂, TiO₂:N, TiO₂:W, and TiO₂:WN films (from bottom-to-top) as a
4 function of each doping scheme. Ti 3p signatures obscured in the case of W containing films as to
5 prioritise visibility of W 4f levels.
6



1
2 Figure 4: (a) Selected UV-Vis transmittance spectra of the sputter-deposited TiO₂, TiO₂:N,
3 TiO₂:W, and TiO₂:WN films (b) Tauc plots of those films with their corresponding extrapolated E_g
4 values.
5



1
2
3
4

Figure 5: 2D-map displaying the variation of E_g as a function of both Nitrogen and Tungsten doping concentrations. The color scale (on the right-hand side) depicts the bandgap values.

ACCEPTED MANUSCRIPT

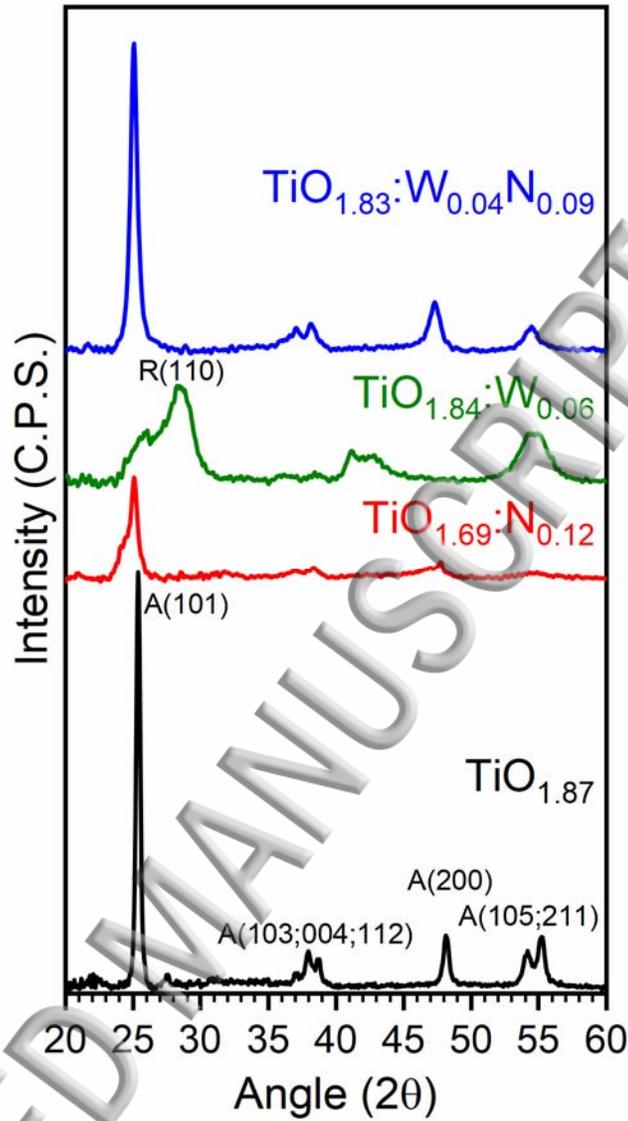
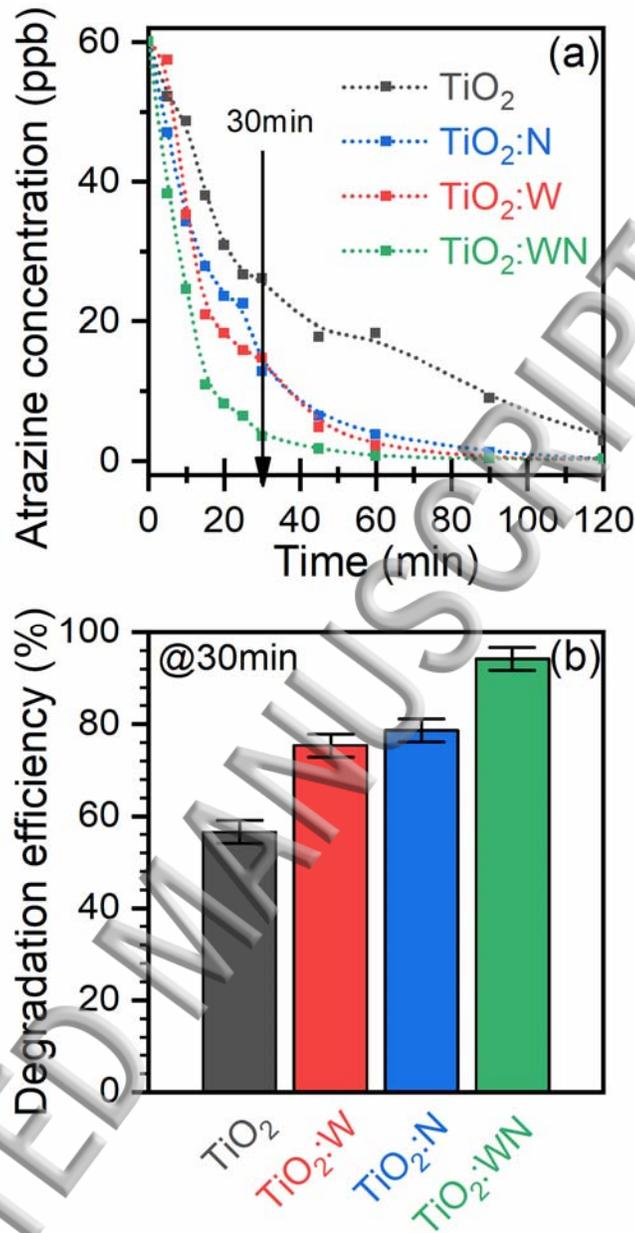


Figure 6: Typical XRD spectra of the sputter-deposited TiO_2 , $\text{TiO}_2\text{:N}$, $\text{TiO}_2\text{:W}$, and $\text{TiO}_2\text{:WN}$ films.

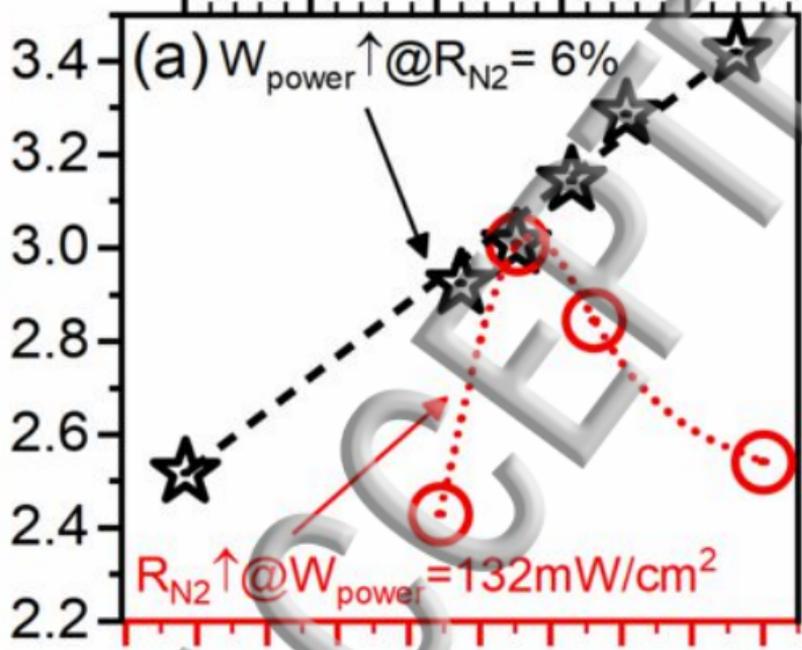


1
2 Figure 7: (a) Residual Atrazine concentration as a function of EPC degradation time, and (b)
3 electro-photocatalytic degradation efficiency of atrazine, under AM1.5 solar irradiation, for the
4 four different photoanodes (i.e. TiO₂, TiO₂:N, TiO₂:W, and TiO₂:WN thin-films).
5

Tungsten Power W_{power} (mW/cm^2)

0 50 100 150 200

Deposition Rate (nm/min)



1 2 3 4 5 6 7 8 9 10

Nitrogen Flow Rate Ratio R_{N_2} (%)

(b)

